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Diffractionless Propagation of Electron Waves in Graphene Superlattices

David E. Fernandes, Manuel Rodrigues, Gabriel Falcão
and Mário G. Silveirinha

*Universidade de Coimbra - Instituto de Telecomunicações
Departamento de Engenharia Electrotécnica e de Computadores*



instituto de
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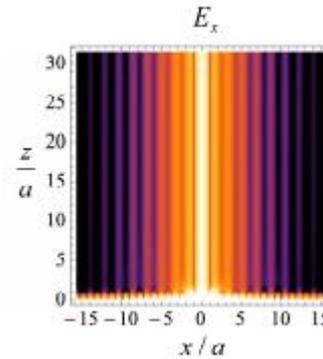
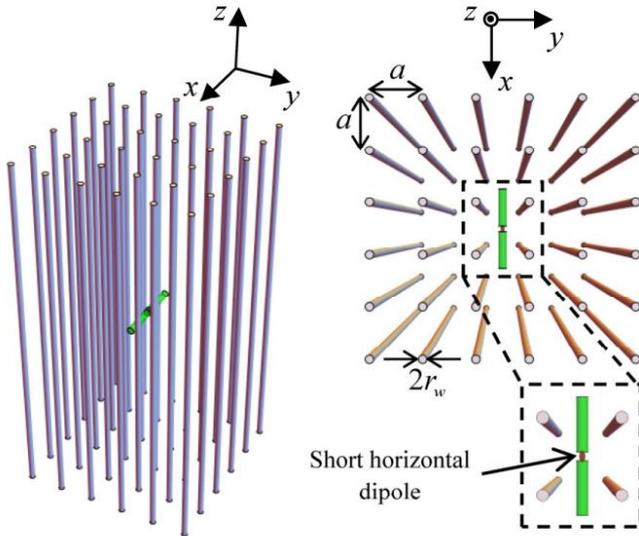
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de leiria

Optical anisotropy

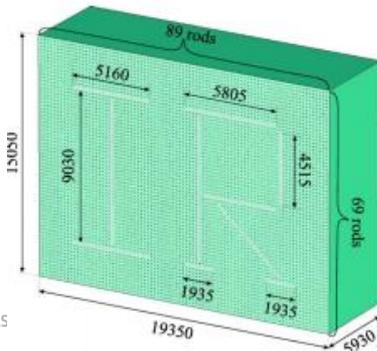
Super collimation of the radiation



Increased Directionality!

T. A. Morgado, M. G. Silveirinha,
<http://arxiv.org/abs/1511.06714>

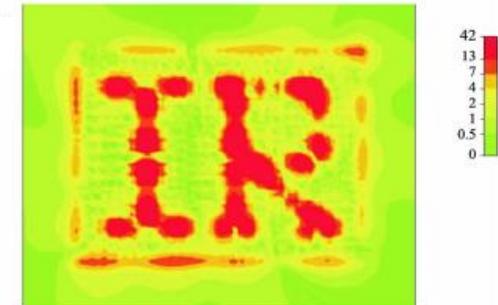
Sub-wavelength imaging



Source Plane



Image Plane



Mário G. Silveirinha, Pavel A. Belov, and Constantin R. Simovski, *Phys. Rev. B* **75**, 035108, 2007
dfernandes@co.it.pt

Electronic anisotropy

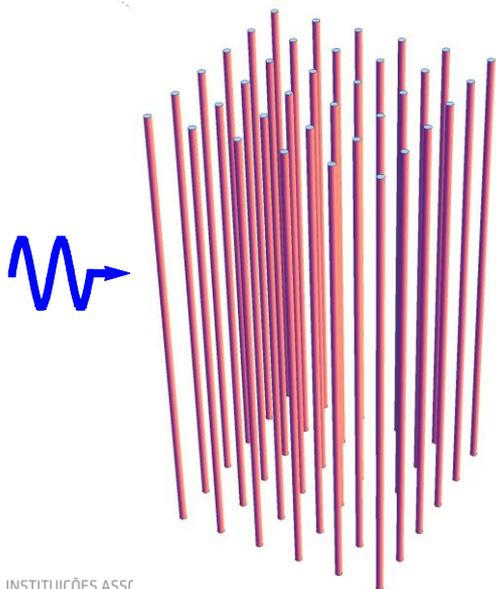
Maxwell Equations

$$\frac{\partial \mathbf{B}(\mathbf{r}, t)}{\partial t} + \nabla \times \mathbf{E}(\mathbf{r}, t) = 0$$

$$\epsilon_0 \frac{\partial \mathbf{E}(\mathbf{r}, t)}{\partial t} + \nabla \times \mathbf{H}(\mathbf{r}, t) = \mathbf{J}_t(\mathbf{r}, t)$$

Schrödinger Equation

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V \right] \psi(\mathbf{r}, t) = i\hbar \frac{\partial \psi(\mathbf{r}, t)}{\partial t}$$



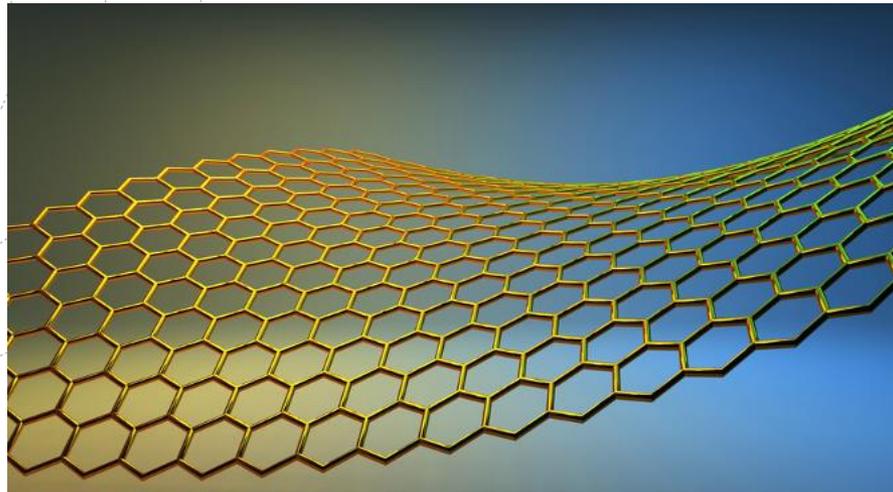
Can we obtain the same phenomena for electrons in graphene based-nanostructures?



<http://tinyurl.com/pbnhz92>

Graphene

Graphene: a two dimensional carbon based material!



<http://tinyurl.com/pbpcuku>



The Nobel Prize in Physics 2010 was awarded jointly to Andre Geim and Konstantin Novoselov *"for groundbreaking experiments regarding the two-dimensional material graphene"*

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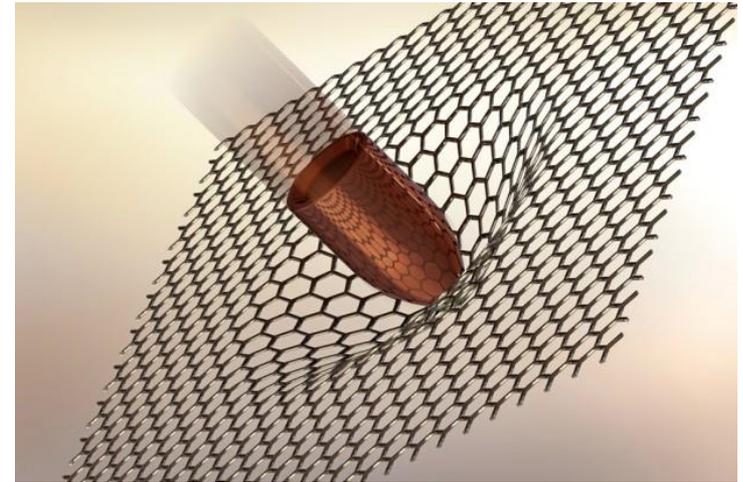


Special Properties of Graphene

- **Flexible and light, yet stronger than steel!!**

Dynamic mechanical behavior of multilayer graphene via supersonic projectile penetration

Jae-Hwang Lee^{1,2,4}, Phillip E. Loya¹, Jun Lou¹, Edwin L. Thomas^{1,4}

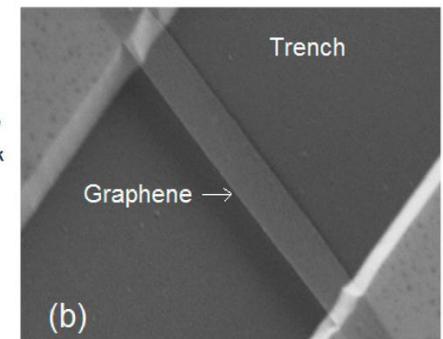
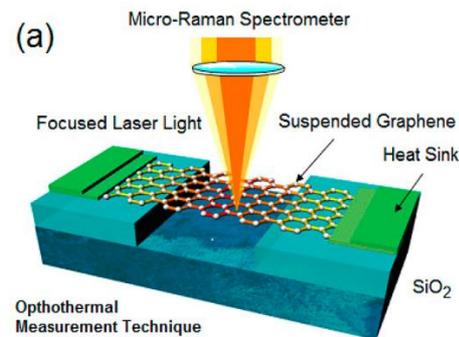


<http://tinyurl.com/oy6zovh>

- **Excellent thermal conductor**

Graphene Thermal Properties: Applications in Thermal Management and Energy Storage

Jackie D. Renteria^{1,2}, Denis L. Nika^{1,3} and Alexander A. Balandin^{1,2,*}

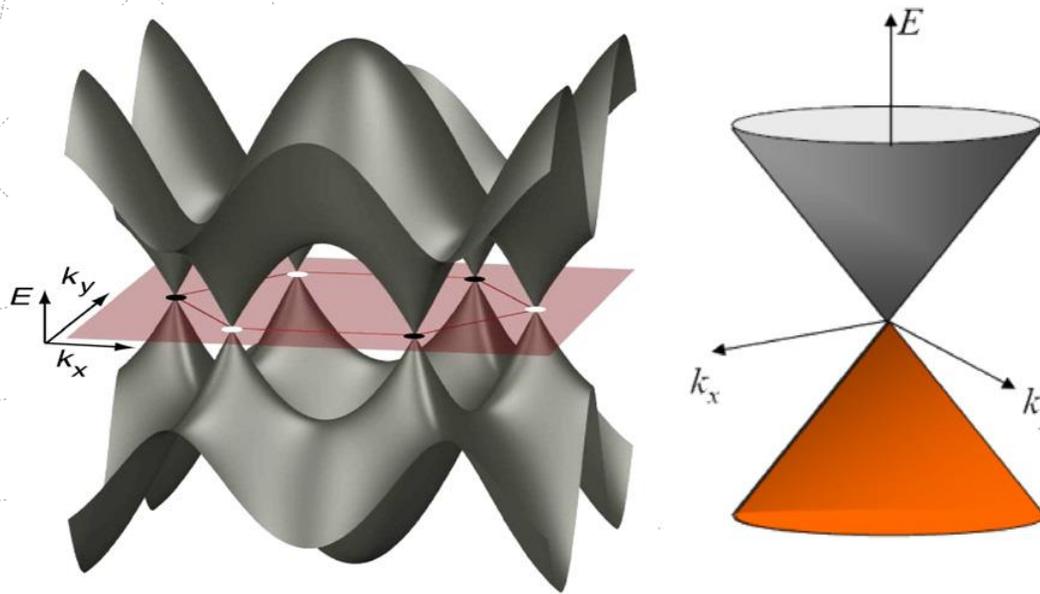


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Electronic Properties of Graphene



Eva Y Andrei *et al*, *Rep. Prog. Phys.* **75** 056501, 2012

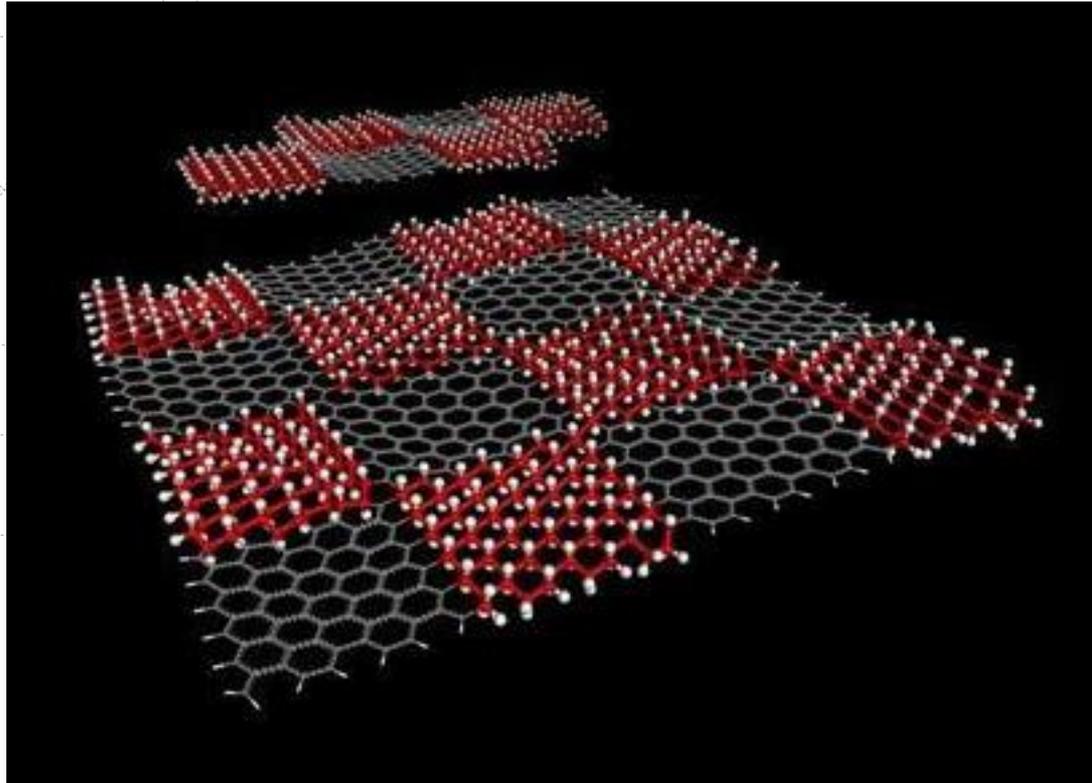
- **Linear energy dispersion.**
- **Zero effective mass.**
- **No bandgap.**

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Graphene Superlattices



<http://goo.gl/JGbgeC>

- **Graphene based heterostructures**
- **Tailoring the potential \leftrightarrow gain control over the electron propagation.**

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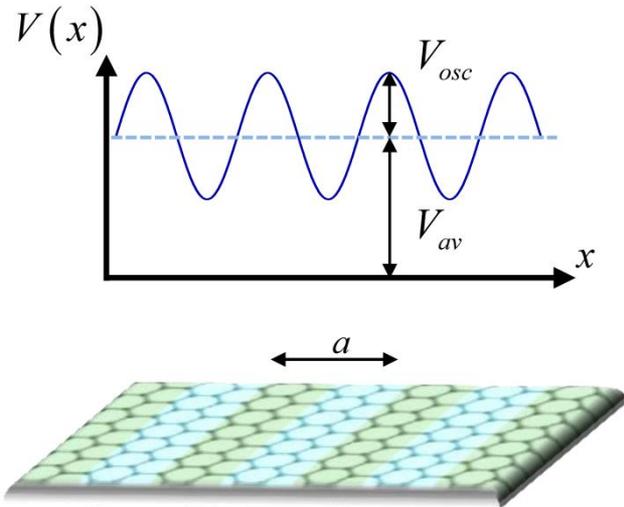


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Graphene Superlattices – Effective Medium Model

Periodic 1-D potential $V(\mathbf{r}) = V(x) = V_{av} + V_{osc}(x)$



$$(\hat{H}\psi)(\mathbf{r}) = [-i\hbar v_F \boldsymbol{\sigma} \cdot \nabla + V(\mathbf{r})] \cdot \psi(\mathbf{r})$$

Effective medium theory

M. G. Silveirinha and N. Engheta
2012 *Phys. Rev. B* **85** 195413

D. E. Fernandes, M. G. Silveirinha and N. Engheta
2014 *Phys. Rev. B* **90**, 041406(R)

$$(\hat{H}_{ef}\psi)(\mathbf{r}) = [-i\hbar v_F (\boldsymbol{\sigma}_x \hat{\mathbf{x}} + \chi \boldsymbol{\sigma}_y \hat{\mathbf{y}}) \cdot \nabla + V_{ef}] \cdot \psi(\mathbf{r})$$

χ anisotropy ratio (depends on $V_{osc}(x)$)

Preferred direction of propagation for the electrons



χ

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Tailoring Transport Properties of Electrons

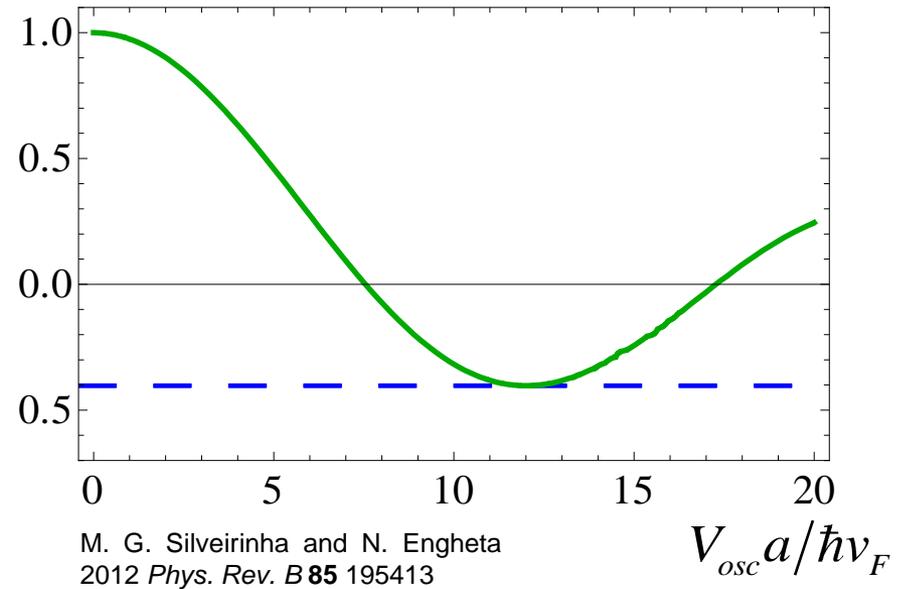
- Anisotropic velocity of propagation

$$v_x = v_F$$

$$v_y = |\chi| v_F$$

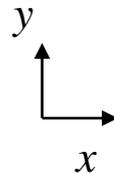
$$V_{osc} \sin\left(\frac{2\pi x}{a}\right)$$

χ

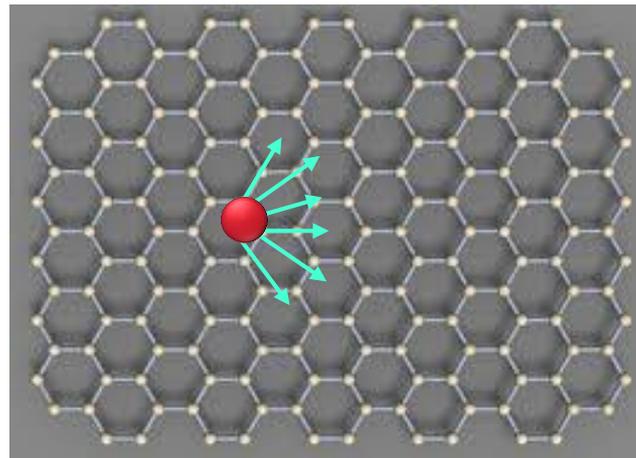


Pristine Graphene: $\chi = 1$

$$v_y = v_x = v_F$$



No preferred direction of propagation!



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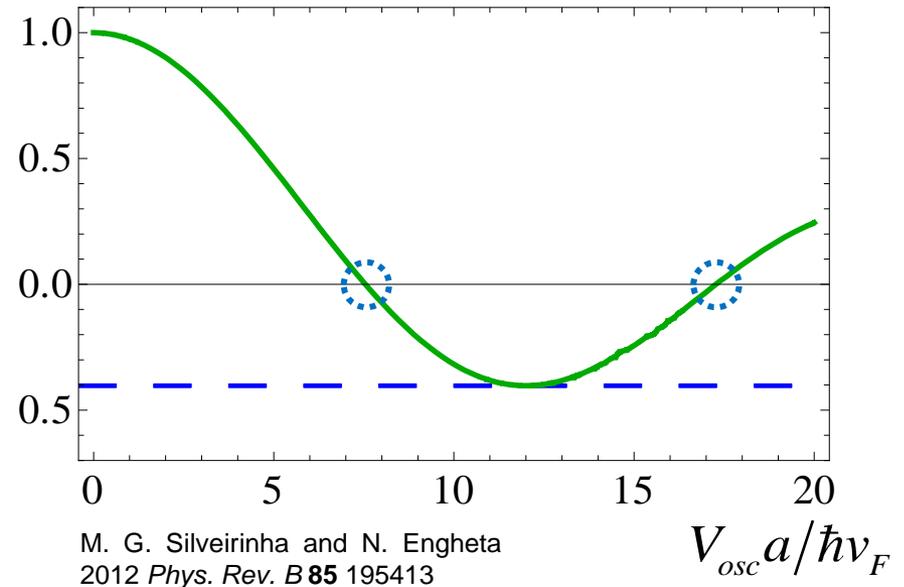
Tailoring Transport Properties of Electrons

- Anisotropic velocity of propagation

$$v_x = v_F$$

$$v_y = |\chi| v_F$$

$$V_{osc} \sin\left(\frac{2\pi x}{a}\right)$$

 χ


Regimes of extreme anisotropy:

$$\frac{V_{osc} a}{\hbar v_F} = 7.55$$

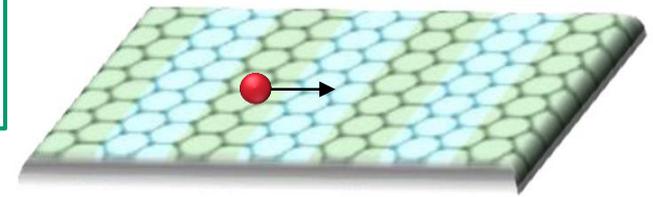
$$\frac{V_{osc} a}{\hbar v_F} = 17.32$$

$$\chi = 0$$

$$v_x = v_F$$

$$v_y = 0$$

Electrons tend to propagate only along the x-direction!!



$x \rightarrow$

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Bare GSL

GSL- Effective medium model

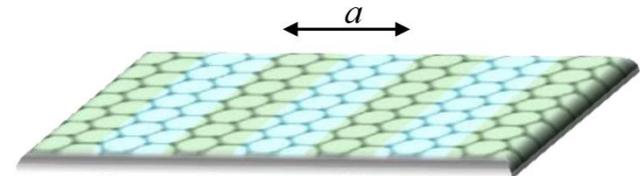
$$\chi = 0$$



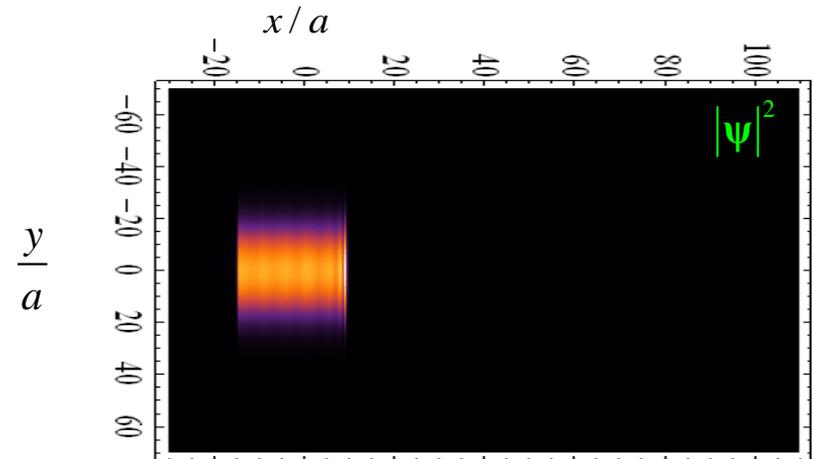
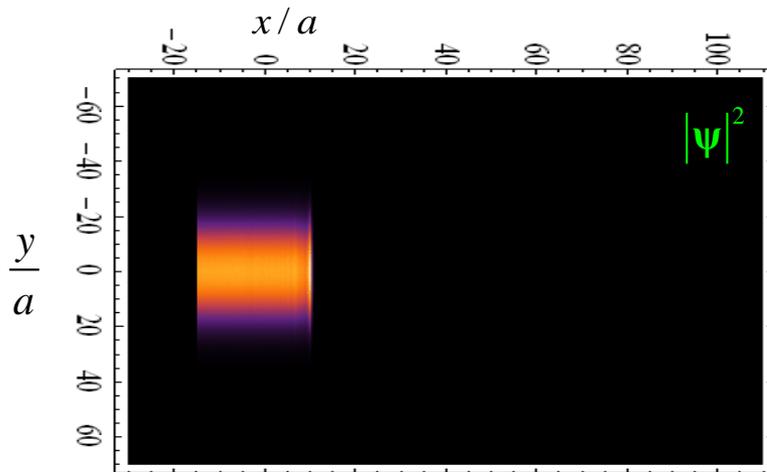
$$R = 10a \quad a = 10\text{nm} \quad (V_{av} - E_0)a/\hbar v_F = 0.1 \quad E_0 a/\hbar v_F = 0.2$$

GSL- Microscopic model

$$\frac{V_{osc} a}{\hbar v_F} = 7.55$$



Time animations



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Effective medium model

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Microscopic model

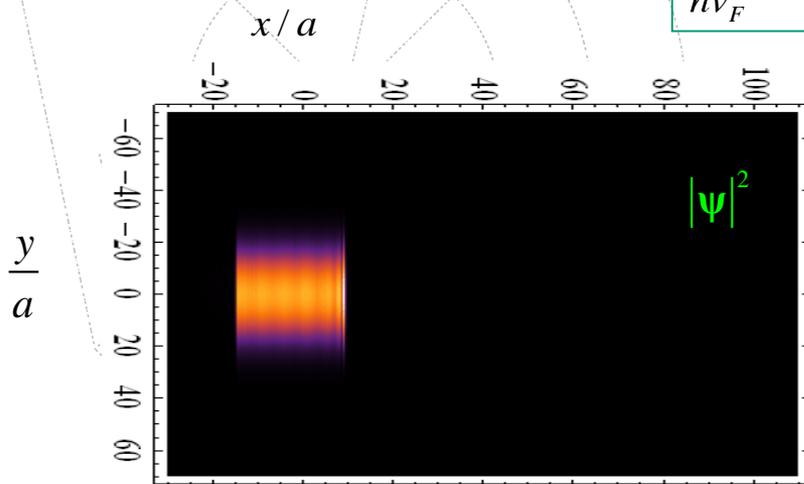


GSLs vs. Pristine Graphene

Time animations

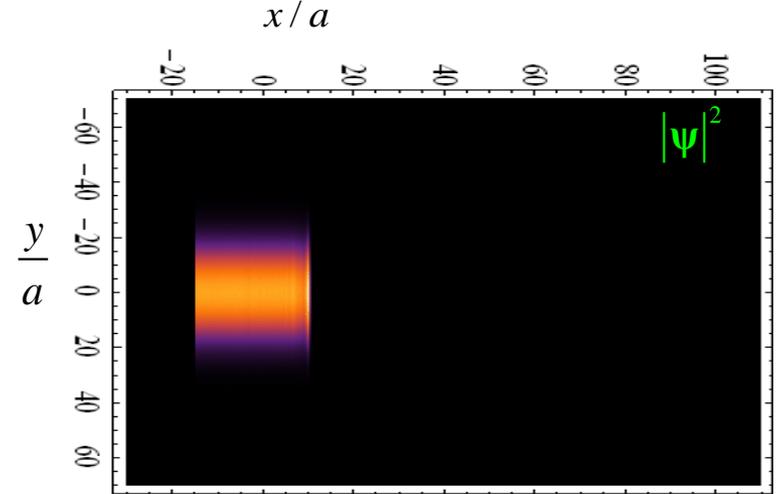
Graphene superlattice – Microscopic model

$$\frac{V_{osc} a}{\hbar v_F} = 7.546$$

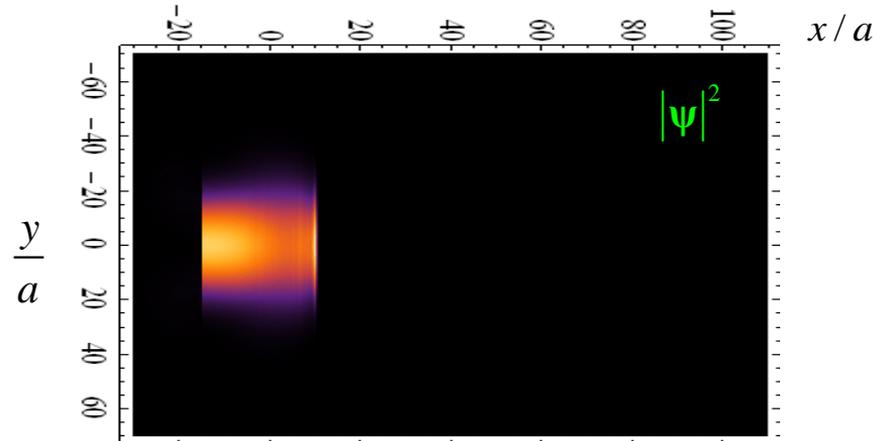


Graphene superlattice – Effective medium model

$$\chi = 0$$



Pristine Graphene $\chi = 1$



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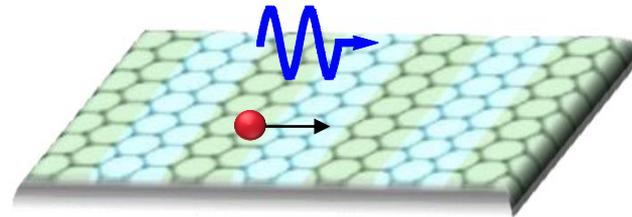


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Conclusions

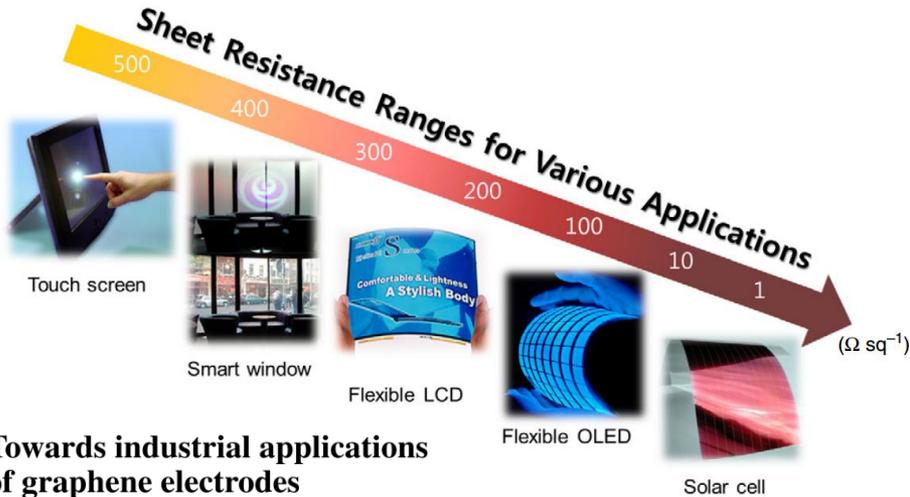
- Tailoring the electric potential in GSLs allows to obtain nondiffractive propagation.
- GSLs may be used for miniaturization of electronic devices.
- Electronic anisotropy creating optical anisotropy!! (Work in progress).



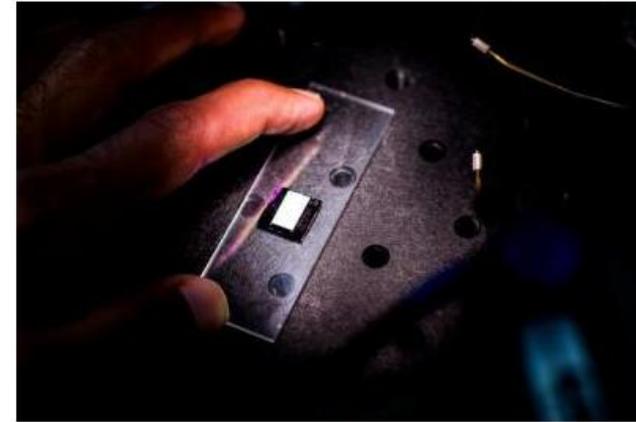
- Enhanced nonlinearities in graphene superlattices (solitons, bistability, high-order harmonics generation ,...).

Potential Applications of Graphene Structures

Transparent and flexible electrodes



Sensors



UM wearable vapor sensor

<http://tinyurl.com/odb7ksz>

Graphene batteries and supercapacitors

Optoelectronics

And much more...

Towards industrial applications of graphene electrodes

Sukang Bae¹, Sang Jin Kim¹, Dolly Shin^{1,3}, Jong-Hyun Ahn^{1,2} and Byung Hee Hong^{1,3}

Solar cells



<http://tinyurl.com/pwf6fvt>

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Energy Dispersion of Graphene Superlattices

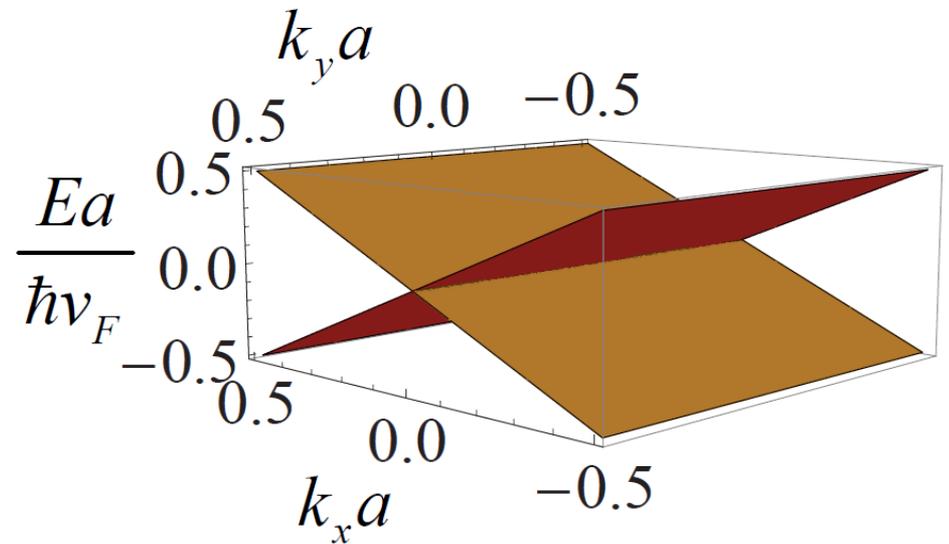
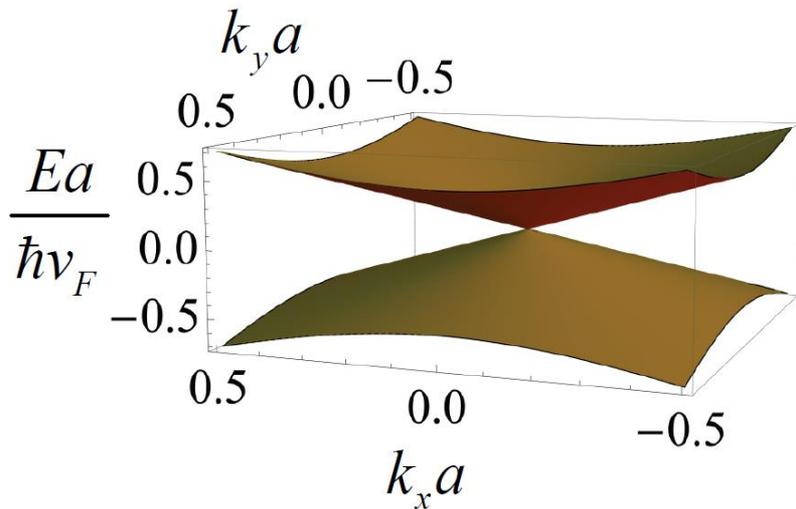
Stationary states energy dispersion:

$$|E| = \hbar v_F \sqrt{k_x^2 + \chi^2 k_y^2}$$

$$\mathbf{v} = \nabla_{\mathbf{k}} E / \hbar = v_F \frac{(k_x, \chi^2 k_y)}{|\mathbf{k}|}$$

Pristine Graphene $\chi = 1$

Superlattice in extreme anisotropy regime $\chi = 0$



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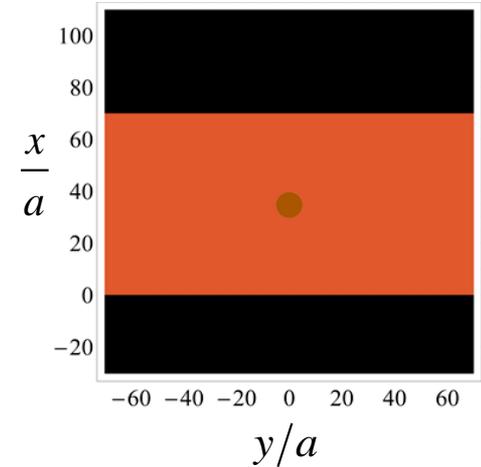
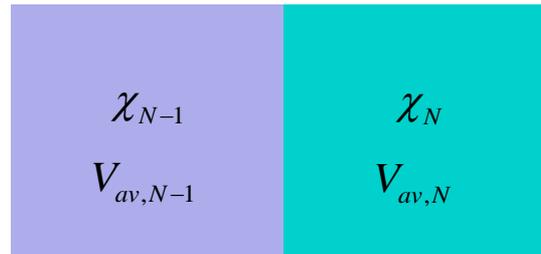
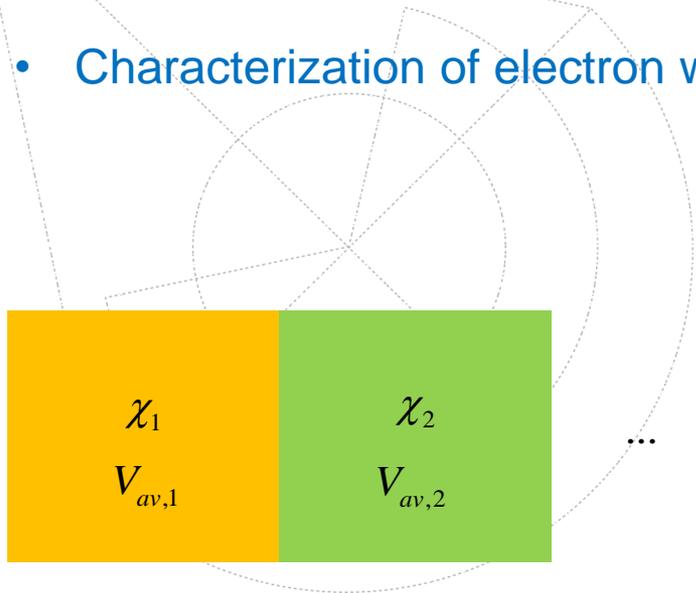


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FDTD algorithm

- Characterization of electron wave propagation in complex GSL structures



- Determining the scattering properties of GSL slabs.
- Characterization of the optical properties of graphene structures (conductivity).
Work in progress...

FDTD Algorithm

- Effective medium model

$$(\hat{H}_{ef} \Psi)(\mathbf{r}) = i\hbar \frac{\partial}{\partial t} \Psi - i\hbar v_F \mathbf{j}$$

$$\hat{H}_{ef} = \left[-i\hbar v_F (\boldsymbol{\sigma}_x \hat{\mathbf{x}} + \chi \boldsymbol{\sigma}_y \hat{\mathbf{y}}) \cdot \nabla + V_{ef} \right]$$

$$\Psi = \begin{Bmatrix} \Psi_1 \\ \Psi_2 \end{Bmatrix}$$

$$\mathbf{j} = \begin{Bmatrix} j_1 \\ j_2 \end{Bmatrix}$$

Fictitious
electron
source

- Time update equations:

$$\frac{\partial \Psi_1}{\partial t} = -v_F \left(\frac{d}{dx} - i\chi \frac{d}{dy} \right) \Psi_2 + \frac{V_{ef}}{i\hbar} \Psi_1 + j_1$$

$$\frac{\partial \Psi_2}{\partial t} = -v_F \left(\frac{d}{dx} + i\chi \frac{d}{dy} \right) \Psi_1 + \frac{V_{ef}}{i\hbar} \Psi_2 + j_2$$

Coupling between

Ψ_1 and Ψ_2

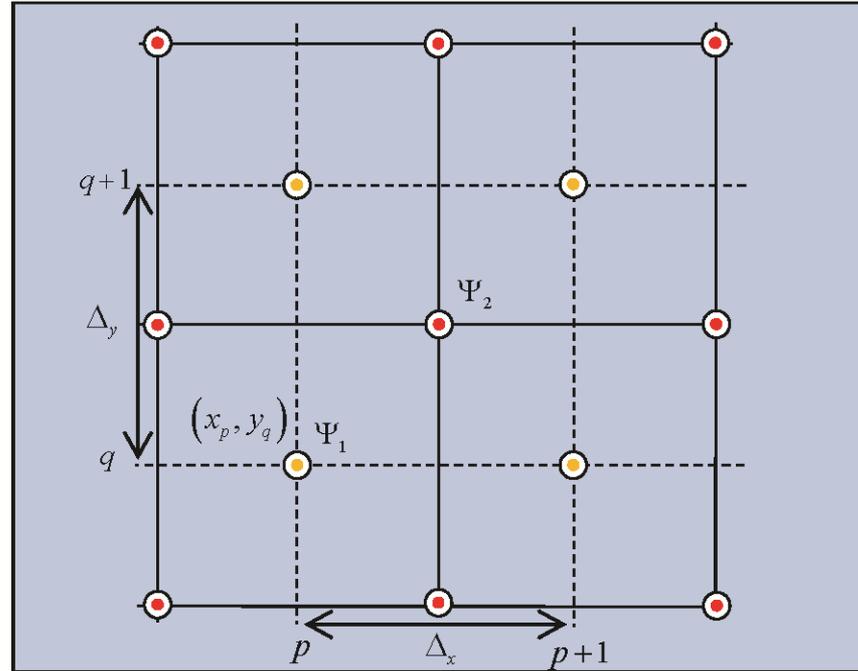
Discretization

- Rectangular Yee-like mesh:

$$F(x, y, t) = F(p\Delta x, q\Delta y, n\Delta t) \equiv F(p, q, n)$$

$$\Psi_1(p, q, n) \equiv \Psi_{1,p,q}^n$$

$$\Psi_2\left(p + \frac{1}{2}, q + \frac{1}{2}, n + \frac{1}{2}\right) \equiv \Psi_{2,p+\frac{1}{2},q+\frac{1}{2}}^{n+\frac{1}{2}}$$



Derivatives → Finite Difference (FD) method

$$\frac{\partial}{\partial i} F(i) = \frac{F(i + \Delta i) - F(i)}{\Delta i}$$

$$i = x, y, t$$



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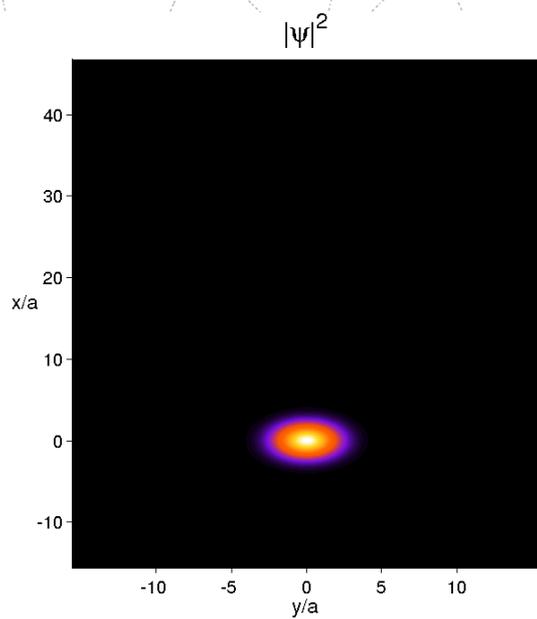


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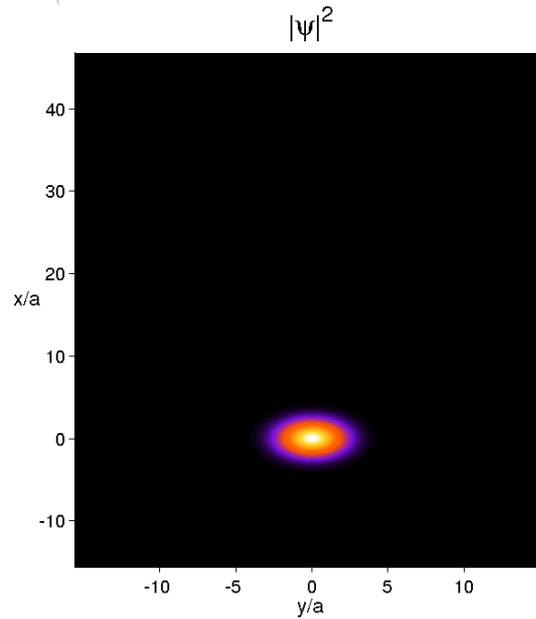
Time Evolution of Initial Electronic States

Time animations



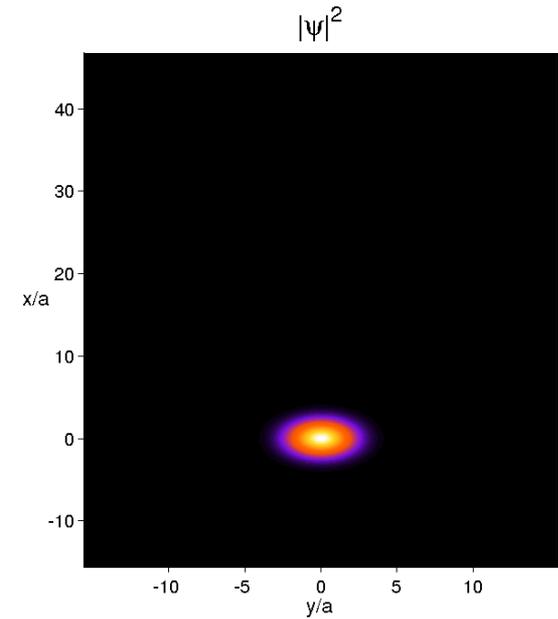
$$\chi = 1.0$$

$$\frac{V_{osc} a}{\hbar v_F} = 0$$



$$\chi = 0.7$$

$$\frac{V_{osc} a}{\hbar v_F} = 3.58$$



$$\chi = 0$$

$$\frac{V_{osc} a}{\hbar v_F} = 7.55$$

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